

Unitary dynamics in strongly Non-Unitary systems



Paolo Zanardi (USC)

AQC 2014 Los Angeles



QIP relies on the ability to generate, preserve & manipulate Q-coherence & Q-entanglement

“Traditional View”

Dissipative dynamics: quantum coherence/entanglement (typically) Destroyed = **Worst enemy of QIP**

Fight with: quantum error correction, decoherence-free Subspaces, dynamical-decoupling, topological-geometric **QIP**,...

“Novel View”

Dissipative dynamics can be *harnessed & exploited* to the end of **QIP**: pure-state preparation, quantum computation, quantum simulations

After all we can make an ally of any enemy, perhaps.....

Setting the Stage

\mathbf{L} = Liouvillian super-operator $\{\mathbf{E}_t\}_{t \geq 0}$ = Evolution semi-group (positive maps)

$$\frac{\partial \mathbf{E}_t}{\partial t} = \mathbf{L} \mathbf{E}_t := (\mathbf{L}_0 + T^{-1} \tilde{\mathbf{K}}) \mathbf{E}_t$$

\mathbf{L}_0 = Dissipative generator $\tilde{\mathbf{K}}(\rho) := -i[\tilde{K}, \rho]$, $\|\tilde{K}\| = O(1)$

Basic Equation: dissipation + small coherent term (T adiabatic time-scale)

Steady-State Manifold (*SSM*): $\rho: \mathbf{L}_0(\rho) = 0$

Q: Can we coherently manipulate states inside the *SSM*?

The Projection Theorem

[PZ & L Campos Venuti, arXiv:1404.4673]

$\mathbf{P}_0 =$ Spectral projection onto $\ker(\mathbf{L}_0)$ ($\mathbf{P}_0\mathbf{L}_0 = \mathbf{L}_0\mathbf{P}_0 = 0$) $\tilde{\mathbf{K}}_{eff} = \mathbf{P}_0\tilde{\mathbf{K}}\mathbf{P}_0$

$$\|(\mathbf{E}_T - e^{\tilde{\mathbf{K}}_{eff}})\mathbf{P}_0\| = O(1/T) \Rightarrow \lim_{T \rightarrow \infty} \mathbf{E}_T\mathbf{P}_0 = e^{\tilde{\mathbf{K}}_{eff}}\mathbf{P}_0$$

$\|(\mathbf{1} - \mathbf{P}_0)\mathbf{E}_T\mathbf{P}_0\| = O(1/T)$ Transitions outside the *SSM* are suppressed

Non SS are adiabatically decoupled and the dynamics inside the *SSM* is ruled by the effective projected generator $\tilde{\mathbf{K}}_{eff}$

- Fermi-Golden rule at infinite order in the perturbation K

Who's $\tilde{\mathbf{K}}_{eff}$?

Projection Theorem: a simple argument

1st order Dyson's series

$$\mathbf{E}_t \mathbf{P}_0 = \mathbf{P}_0 + \frac{t}{T} \mathbf{P}_0 \tilde{\mathbf{K}} \mathbf{P}_0 + \frac{1}{T} \int_0^t e^{t\mathbf{L}_0} \mathbf{Q}_0 \tilde{\mathbf{K}} \mathbf{P}_0 dt + O(T^{-2})$$

$$\int_0^T e^{t\mathbf{L}_0} \mathbf{Q}_0 \tilde{\mathbf{K}} \mathbf{P}_0 dt = \mathbf{S} (e^{T\mathbf{L}_0} - \mathbf{1}) \tilde{\mathbf{K}} \mathbf{P}_0 \quad \mathbf{S} := - \int_0^\infty e^{t\mathbf{L}_0} \mathbf{Q}_0 dt = \mathbf{L}_0^{-1}$$

$\text{Re}(\lambda_h) < 0 \quad (h > 0)$

$$\| \mathbf{E}_T \mathbf{P}_0 - e^{\mathbf{P}_0 \tilde{\mathbf{K}} \mathbf{P}_0} \mathbf{P}_0 \| \leq 2 \frac{\| \mathbf{S} \|}{T} \| \tilde{\mathbf{K}} \| + O(T^{-2}), \quad \| \mathbf{S} \| = O(\Delta^{-m_h})$$

- Deviations from adiabaticity i.e., transitions outside the **SSM** controlled by dissipative gap $\Delta := \min_{h>0} | \text{Re} \lambda_h |$

- Dynamical phase transitions, info. geometry of diss. **QPTs**

Example 1

$$\mathbf{L}_0(\rho) = \Phi(\rho) - \rho, \quad \Phi(\rho) = \sum A_i \rho A_i^+, \quad \Phi(\mathbf{1}) = \mathbf{1} \quad \Phi^*(\mathbf{1}) = \mathbf{1}$$

$$A := \text{Alg}[A_i, A_i^+] \Rightarrow \ker \mathbf{L}_0 = A' \cong \bigoplus_J M(n_J) \otimes \mathbf{1}_{d_J}$$

Dissipator kernel =
commutant algebra
(Kribs 2006)

SSM: convex hull of A -noiseless states (mixed in general)

$$\mathbf{P}_0(X) = \int_{U \in \mathbf{U}(A)} UXU^+ dU = \sum_J \text{Tr}_{d_J}(\Pi_J X \Pi_J) \otimes \mathbf{1}_{d_J} \in A' = A\text{-symmetrization}$$

$$\tilde{\mathbf{K}}_{\text{eff}}(X) = -i[\tilde{K}_{\text{eff}}, X], \quad \tilde{K}_{\text{eff}} = P_0(\tilde{K}) \quad (X \in A')$$

Effective generator: environment-symmetrized K

Example 2: one free system coupled to a dissipative one

$$\mathbf{L}_0 = \mathbf{1}_A \otimes \mathbf{L}_B, \quad \exists! \rho_B / \mathbf{L}_B(\rho_B) = 0 \quad \Rightarrow \quad \mathbf{L}_0(X \otimes \rho_B) = 0 \quad \forall X$$

$$\tilde{\mathbf{K}}_{eff}(X) = -i[\tilde{K}_{eff}, X], \quad \tilde{K}_{eff} = Tr_B[(\mathbf{1}_A \otimes \rho_B)K] \otimes \mathbf{1}_B$$

Example 3: **SSM** \supset a Decoherence-Free Subspace

$$|\phi\rangle, |\phi'\rangle \in C \subset H, \quad |\varphi\rangle \in C^{perp} \quad \Rightarrow \quad \mathbf{L}_0(|\phi\rangle\langle\phi'|) = 0 \quad \mathbf{P}_0(|\phi\rangle\langle\varphi|) = \mathbf{P}_0(|\varphi\rangle\langle\phi|) = 0$$

$$X \in L(C) \quad \Rightarrow \quad \tilde{\mathbf{K}}_{eff}(X) = -i[\tilde{K}_{eff}, X] \quad \tilde{K}_{eff} = \Pi_C \tilde{K} \Pi_C$$

Effective dynamics over the **DFS** ruled by a projected Hamiltonian

Remarks I

- For **Abelian** A : $\mathbf{P}_0(X) = \sum_J \Pi_J X \Pi_J =$ meas. super-op: **Zeno dynamics**

G. A. Paz-Silva et PRL 108, 080501 (2012); D. Burgath et al, arXiv:1403.5752

- Rotating frame generated by K : no Hamiltonian term & (slowly) time-dependent bath: The **Projection Theorem is an adiabatic th. for open Sys** (i.e. system's always inside the instantaneous SSM)

M. S. Sarandy, D. A. Lidar, PRA 73, 062101 (2006)

- Effective evolution in the SSM is geometric → **dissipation-assisted Holonomic Quantum Computation**

A. Carollo et PRL 96, 020403 (2006); O. Oreshkov, J. Calsamiglia, PRL, 105 050503 (2010)

Remarks II

- Effective dynamics may be *more complex* than the bare one:
More-non local, higher-computational power e.g.,

$$[K_1, K_2] = 0 \text{ but } [P_0(K_1), P_0(K_2)] \neq 0$$

D. Burgath et al, arXiv:1403.5752

- Effective dynamics is *robust against perturbations* of L
that are removed by the projection e.g.,

$$H - K \in \text{Ker}(P_0) \implies H_{\text{eff}} = K_{\text{eff}}$$

- If the 1st order effective terms vanish one has a 2nd order effective
(non unitary) dynamics over the **SSM** (up to $O\left(\frac{\|S\|}{T^{1/2}}\right)$ corrections)

$$P_0 \tilde{K} P_0 = 0, \quad \|K\| = O(T^{-1/2}) \quad \implies \quad \tilde{K}_{\text{eff}} = - P_0 \tilde{K} S \tilde{K} P_0$$

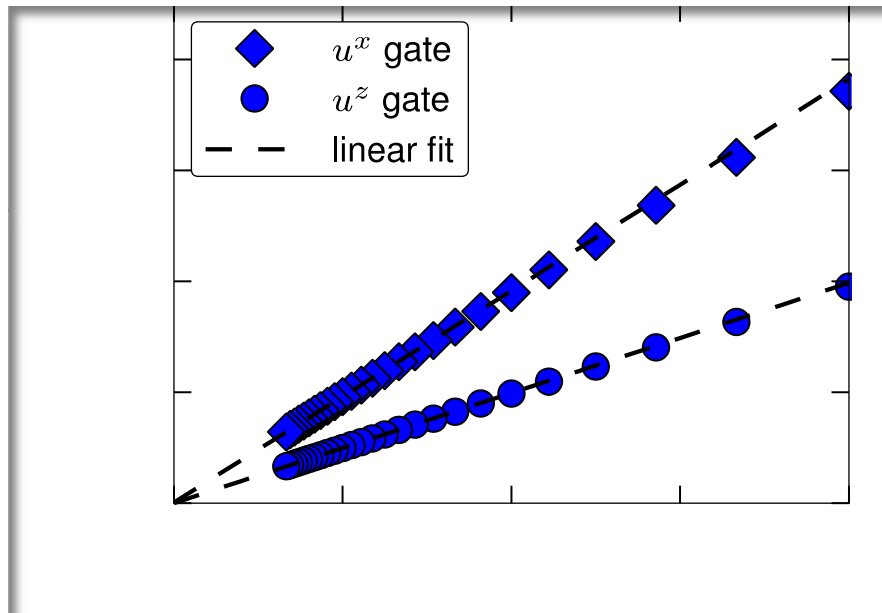
Dissipation-assisted logic gates on a DFS

$$\mathbf{L}_0(\rho) = \sum_{\alpha=x,y,z} (S^\alpha \rho S^\alpha - 1/2 \{(S^\alpha)^2, \rho\}), \quad S^\alpha := \sum_{j=1}^N \sigma_j^\alpha$$

$$N=4 \implies \ker(\mathbf{L}_0) = M(2) \otimes \mathbf{1}_1 \oplus \dots \cong L(C) \oplus \dots \quad [\dim(\ker \mathbf{L}_0) = 14]$$

$$C := \{|\phi\rangle \in (\mathbb{C}^2)^{\otimes 4} / S^\alpha |\phi\rangle = 0, \alpha = x, y, z\}$$

Singlet sector = *DFS*



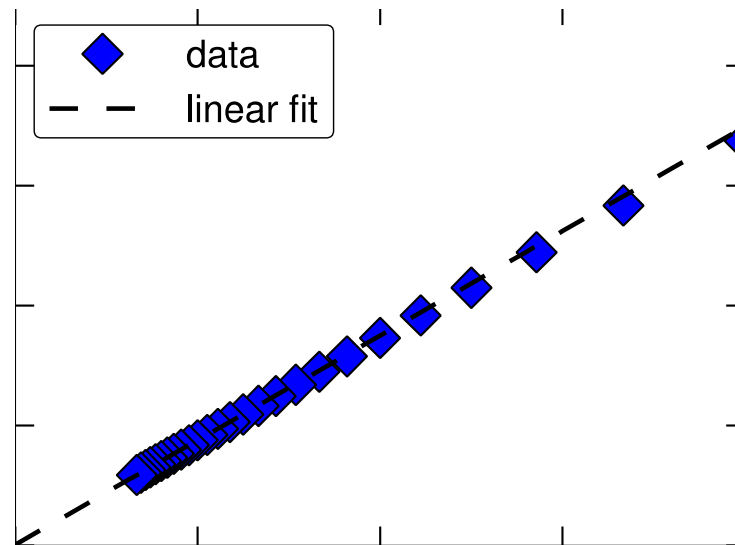
$$\tilde{K}^x = \frac{3}{2}(\sigma_1^z \sigma_2^z + \sigma_2^z \sigma_3^z + \mathbf{1}), \quad \tilde{K}^z = -\frac{\sqrt{3}}{2}(\sigma_1^z \sigma_2^z - \sigma_2^z \sigma_3^z + \sigma_1^z), \quad \implies \tilde{K}_{eff}^\alpha \approx \sigma^\alpha$$

Dissipation-assisted logic gates on Noiseless Subsystems



$$\mathbf{L}_0(\rho) = 1/3 \sum_{\alpha=x,y,z} U^\alpha \rho (U^\alpha)^\dagger - \rho, \quad U_\alpha := e^{i\theta_\alpha S^\alpha}$$

$$N=3 \quad \ker \mathbf{L}_0 = M(2) \otimes \mathbf{1}_2 \oplus M(1) \otimes \mathbf{1}_4 \quad \dim(\ker \mathbf{L}) = 2^2 + 1^1 = 5$$



$$\tilde{K} = \sigma_1^x \sigma_2^x$$

Remark: The *SSM* comprises just mixed states

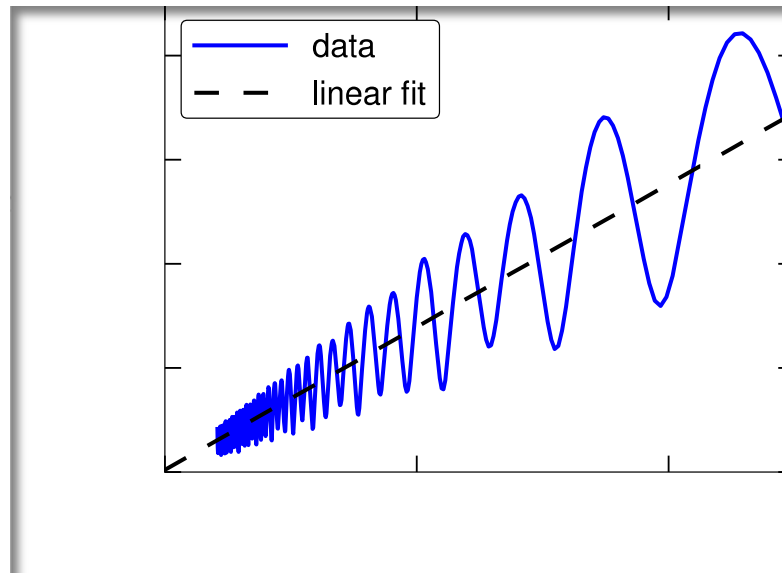
A Unitary example

$$H_0 = S^+ \otimes \left(\sum_q g_q b_q \right) + S^- \otimes \left(\sum_q g_q b_q^+ \right) + \mathbf{1} \otimes \sum_q \omega_q b_q^+ b_q$$

$$|\phi_k\rangle := \sum_{j=1}^N e^{i \frac{2\pi j k}{N}} \sigma_j^+ |\downarrow, \dots, \downarrow\rangle \otimes |0\rangle \Rightarrow H_0 |\phi_k\rangle = 0 \quad (k=1, \dots, N-1)$$

$$C = \text{span} \{ |\phi_k\rangle \}_{k=1}^{N-1}$$

Dark state manifold of H_0 (DFS @ $T=0$)



$$\tilde{K} = \sigma_1^z \Rightarrow \tilde{K}_{\text{eff}} = 2(N-1)|\phi\rangle\langle\phi| - \mathbf{1}, \quad |\phi\rangle := \frac{1}{\sqrt{N-1}} \sum_{k=1}^{N-1} e^{-i \frac{2\pi k}{N}} |\phi_k\rangle$$

Projection Th reduces to an adiabatic one for closed sys

Conclusions

- Projection th. over the *SSM* of strongly dissipative systems
- Adiabatic coherent (geometric) manipulations on the *SSM*
- Complexity, (non) locality and robustness can be enhanced
- Dissipation-assisted gates in *DFSs* and *NSs*

***Moral:** dissipation & decoherence can be good guys...*

PZ & L Campos Venuti, arXiv:1404.4673

THANKS FOR THE ATTENTION!

Projection Theorem: Sketch of the proof

\mathbf{P} spectral projection of \mathbf{L} originated by the 0 eigenvalue of \mathbf{L}_0

$$\exp(t\mathbf{L})\mathbf{P} = \exp(t\mathbf{P}\mathbf{L}\mathbf{P})\mathbf{P} = \exp[t\mathbf{P}_0\mathbf{L}\mathbf{P}_0 + t(\mathbf{P}\mathbf{L}\mathbf{P} - \mathbf{P}_0\mathbf{L}\mathbf{P}_0)]\mathbf{P} := e^{\mathbf{X}+\mathbf{Y}}\mathbf{P}$$

$$\mathbf{X} := t\mathbf{P}_0\mathbf{L}\mathbf{P}_0 = \frac{t}{T}\mathbf{P}_0\tilde{\mathbf{K}}\mathbf{P}_0$$

$$\mathbf{Y} := t(\mathbf{P}\mathbf{L}\mathbf{P} - \mathbf{P}_0\mathbf{L}\mathbf{P}_0) = -\frac{t}{T^2}\left(\mathbf{P}_0\tilde{\mathbf{K}}\mathbf{P}_0\tilde{\mathbf{K}}\mathbf{S} + \mathbf{P}_0\tilde{\mathbf{K}}\mathbf{S}\tilde{\mathbf{K}}\mathbf{P}_0 + \mathbf{S}\tilde{\mathbf{K}}\mathbf{P}_0\tilde{\mathbf{K}}\mathbf{P}_0\right) + O\left(\frac{t}{T^3}\right)$$

$$\mathbf{P} = \mathbf{P}_0 + (\mathbf{P} - \mathbf{P}_0) = \mathbf{P}_0 - \frac{1}{T}\left(\mathbf{P}_0\tilde{\mathbf{K}}\mathbf{S} + \mathbf{S}\tilde{\mathbf{K}}\mathbf{P}_0\right) + O\left(\frac{1}{T^2}\right) \quad (\text{Kato})$$

$$\exp(t\mathbf{L})\mathbf{P}_0 = [e^{\mathbf{X}} + (e^{\mathbf{X}+\mathbf{Y}} - e^{\mathbf{X}})]\mathbf{P} - \exp(t\mathbf{L})(\mathbf{P} - \mathbf{P}_0) \Rightarrow$$

$$\exp(t\mathbf{L})\mathbf{P}_0 - e^{\mathbf{X}}\mathbf{P}_0 = e^{\mathbf{X}}(\mathbf{P} - \mathbf{P}_0) + (e^{\mathbf{X}+\mathbf{Y}} - e^{\mathbf{X}})\mathbf{P} - e^{t\mathbf{L}}(\mathbf{P} - \mathbf{P}_0)$$

$$\|\exp(t\mathbf{L})\mathbf{P}_0 - e^{\mathbf{X}}\mathbf{P}_0\| \leq 2\|\mathbf{P} - \mathbf{P}_0\| + \|e^{\mathbf{X}+\mathbf{Y}} - e^{\mathbf{X}}\| \leq 2\|\mathbf{P} - \mathbf{P}_0\| + \|\mathbf{Y}\| e^{\|\mathbf{X}\|+\|\mathbf{Y}\|}$$

$$\Rightarrow \quad \|e^{t\mathbf{L}}\mathbf{P}_0 - e^{\mathbf{P}_0\tilde{\mathbf{K}}\mathbf{P}_0}\mathbf{P}_0\| = O\left(\frac{1}{T}\right) \quad q.e.d.$$

$$\mathbf{L}_0 = \mathbf{1}_A \otimes \mathbf{L}_B, \quad \exists! \rho_B / \mathbf{L}_B(\rho_B) = 0 \Rightarrow \mathbf{L}_0(X \otimes \rho_B) = 0 \quad \forall X$$

$$\tilde{\mathbf{K}}_{eff}(X) = -i[\tilde{K}_{eff}, X], \quad \tilde{K}_{eff} = Tr_B[(\mathbf{1}_A \otimes \rho_B)K] \otimes \mathbf{1}_B$$

$$|\phi\rangle \in C \subset H, \quad |\varphi\rangle \in C^{perp} \Rightarrow \mathbf{L}_0(|\phi\rangle\langle\phi|) = |\phi\rangle\langle\phi| \quad \mathbf{L}_0(|\phi\rangle\langle\varphi|) = \mathbf{L}_0(|\varphi\rangle\langle\phi|) = 0$$

$$X \in L(C) \Rightarrow \tilde{\mathbf{K}}_{eff}(X) = -i[\tilde{K}_{eff}, X] \quad \tilde{K}_{eff} = \Pi_C \tilde{K} \Pi_C$$

$$\mathbf{SSM} = Conv-Hull \left\{ \rho_J \otimes \frac{\mathbf{1}_{d_J}}{d_J} \right\}$$

$$\ker \mathbf{L}_0 \cong \bigoplus_J M(n_J) \otimes \mathbf{1}_{d_J}, \quad n_J = \frac{(2J+1)N!}{(N/2+J+1)!(N/2-1)!},$$

$$\mathbf{P}_0(X) = \int_{U \in \mathbf{U}(A)} UXU^\dagger dU = \sum_J Tr_{d_J}(\Pi_J X \Pi_J) \otimes \mathbf{1}_{d_J} \in A'$$

$$X \in A' \Rightarrow \tilde{\mathbf{K}}_{eff}(X) = -i[\tilde{K}_{eff}, X], \quad \tilde{K}_{eff} = P_0(\tilde{K})$$

$$\mathbf{E}_t \mathbf{P}_0 = \mathbf{P}_0 + \frac{t}{T} \mathbf{P}_0 \tilde{\mathbf{K}} \mathbf{P}_0 + \frac{1}{T} \int_0^t e^{t\mathbf{L}_0} \mathbf{Q}_0 \tilde{\mathbf{K}} \mathbf{P}_0 dt + O(T^{-2})$$

$$\int_0^T e^{t\mathbf{L}_0} \mathbf{Q}_0 \tilde{\mathbf{K}} \mathbf{P}_0 dt = \mathbf{S}(e^{T\mathbf{L}_0} - \mathbf{1}) \tilde{\mathbf{K}} \mathbf{P}_0$$

$$\mathbf{S} := - \int_0^\infty e^{t\mathbf{L}_0} \mathbf{Q}_0 dt = - \sum_{h>0} \left(\frac{\mathbf{P}_h}{(-\lambda_h)} + \sum_{m=1}^{m_h-1} \frac{\mathbf{D}_h^m}{(-\lambda_h)^{m+1}} \right) = \text{"}\mathbf{L}_0^{-1}\text{"}$$

$$\| \mathbf{E}_T \mathbf{P}_0 - e^{\mathbf{P}_0 \tilde{\mathbf{K}} \mathbf{P}_0} \mathbf{P}_0 \| \leq 2 \frac{\| \mathbf{S} \|}{T} \| \tilde{\mathbf{K}} \| + O(T^{-2})$$

$$\| \mathbf{S} \| = O(\Delta^{-m_h}) \quad \Delta := \min_{h>0} | \lambda_h |$$

$$\mathbf{R}(z) := (z - \mathbf{L})^{-1} = \sum_h \left(\frac{\mathbf{P}_h}{z - \lambda_h} + \sum_{k=1}^{m_h-1} \frac{\mathbf{D}_h^k}{(z - \lambda_h)^{k+1}} \right) \quad \mathbf{P}_0 \mathbf{L}_0 = \mathbf{L}_0 \mathbf{P}_0 = \mathbf{0}$$

$$\mathbf{L} = \sum_h (\lambda_h \mathbf{P}_h + \mathbf{D}_h), \quad \mathbf{P}_h \mathbf{D}_h = \mathbf{D}_h \mathbf{P}_h = \mathbf{D}_h, \quad [\mathbf{P}_h, \mathbf{P}_k] = 0, \quad \sum_h \mathbf{P}_h = \mathbf{1}, \quad (\mathbf{D}_h)^{m_h} = \mathbf{0}$$

$$\frac{\partial \mathbf{E}_t}{\partial t} = \mathbf{L} \mathbf{E}_t := (\mathbf{L}_0 + T^{-1} \tilde{\mathbf{K}}) \mathbf{E}_t \Rightarrow \mathbf{E}_t = e^{t\mathbf{L}}$$

$$\tilde{\mathbf{K}}(\rho) := -i[\tilde{K}, \rho], \quad \|\tilde{K}\| = O(1)$$

$$\|(\mathbf{E}_T - e^{\tilde{\mathbf{K}}_{eff}}) \mathbf{P}_0\| = O(1/T) \quad \tilde{\mathbf{K}}_{eff} := \mathbf{P}_0 \tilde{\mathbf{K}} \mathbf{P}_0$$

$$\mathbf{E}_t = e^{t\mathbf{L}} = \prod_h e^{t(\lambda_h \mathbf{P}_h + \mathbf{D}_h)} = \mathbf{P}_0 + \sum_{h>0} (\mathbf{P}_h e^{t\lambda_h} + e^{t\lambda_h} \sum_{k=1}^{m_h-1} \frac{(t\mathbf{D})^k}{k!})$$